# A Schiffer-type problem for annular domains joint work with A. Enciso, A. J. Fernández and P. Sicbaldi 

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IMAG, University of Granada
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## Outline

Introduction: the Pompeiu problem and the Schiffer conjecture

A Schiffer-type problem

Proof of the main result

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Introduction: the Pompeiu problem and the Schiffer conjecture

## A Schiffer-type problem

## Proof of the main result

## The Pompeiu problem

Let $f: \mathbb{R}^{N} \rightarrow \mathbb{R}$ be a continuous function and $\Omega \subset \mathbb{R}^{N}$ a bounded domain. The Pompeiu problem consists in recovering $f$ from the values:

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where $\mathcal{R}$ is any rigid motion in $\mathbb{R}^{N}$.
We say that $\Omega$ satisfies the Pompeiu property if $f=0$ is the unique function $f$

$$
\int_{\mathcal{R}(\Omega)} f(x) d x=0
$$

for any rigid motion $\mathcal{R}$.

The Pompeiu problem


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The Pompeiu problem


## The Pompeiu problem



Ellipses and polygons satisfy the Pompeiu property (Brown, Schreiber \& Taylor '73).

## Balls do not satisfy the Pompeiu property!

Let $\Omega$ be any ball of radius R , and $\mu \neq 0$ a radial Neumann eigenvalue; then,

$$
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where $w(x)=(\mu c)^{-1}(u(x)-c)$.
We now take $f: \mathbb{R}^{2} \rightarrow \mathbb{R}$ such that $-\Delta f=\mu f$. For instance, $f(x)=\sin \left(\sqrt{\mu} x_{1}\right)$, or $f(x)=J_{\frac{N-2}{2}}(\sqrt{\mu}|x|)$.

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## An overdetermined elliptic problem

Observe that the above argument works as long as there exists a solution to the problem:

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Theorem (Williams '76)
If $\partial \Omega$ is homeomorphic to $\mathbb{S}^{N-1}$ and $\Omega$ does not satisfy the Pompeiu property, then there exists a solution to ( S ).

The topological assumption is necessary!

## The Schiffer conjecture

Schiffer Conjecture (problem 80 in Yau's list)
If problem ( S ) admits a solution, then $\Omega$ is a ball.

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The Schiffer conjecture is open, but there are some results related to it. We state below some of them:

1. If $\Omega$ is $C^{1}$ and there exists a solution to ( S ), then $\Omega$ is analytic (Kinderlehrer-Nirenberg '77).

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1. If $\Omega$ is $C^{1}$ and there exists a solution to ( S ), then $\Omega$ is analytic (Kinderlehrer-Nirenberg '77).
2. If $\partial \Omega$ is connected and there exists a diverging sequence $\mu_{k}$ for which (S) is solvable, then $\Omega$ is a ball (Berenstein '80, Berenstein-Yang '87).

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1. If $\Omega$ is $C^{1}$ and there exists a solution to ( S ), then $\Omega$ is analytic (Kinderlehrer-Nirenberg '77).
2. If $\partial \Omega$ is connected and there exists a diverging sequence $\mu_{k}$ for which ( S ) is solvable, then $\Omega$ is a ball (Berenstein '80, Berenstein-Yang '87).
3. If $N=2, \partial \Omega$ is connected and ( S ) admits a solution for
$\mu \leq \mu_{6}$, then $\Omega$ is a disk (Avilés '86, Deng '12). Here we are denoting the Neumann eigenvalues as:

$$
0=\mu_{0}<\mu_{1} \leq \mu_{2} \ldots
$$

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## A related question

In this talk we are interested in the following question:

## Question

Let $\Omega \subset \mathbb{R}^{2}$ a domain, and let us denote $\Gamma_{i}$ the connected components of $\partial \Omega$. Assume that there exists a solution to the problem:

$$
\left\{\begin{array}{ll}
-\Delta u=\mu u & \text { in } \Omega,  \tag{Q}\\
u=c_{i} & \text { on } \Gamma_{i}, \\
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Is it true that $\Omega$ is a disk or an annulus?

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Is it true that $\Omega$ is a disk or an annulus?
Observe that here, as in problem (S), we have that $\nabla u=0$ on $\partial \Omega$.

## Rigidity results

This problem shares many features with the Schiffer conjecture. To start with, if $\Omega$ is $C^{1}$ and there exists a solution to $(Q)$, then $\Omega$ is analytic (Kinderlehrer-Nirenberg '77).

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## Proposition

If there exists a diverging sequence $\mu_{k}$ for which ( Q ) is solvable, then $\Omega$ is a disk or an annulus.
The proof uses the same ideas of Berenstein together with a unique continuation argument.

## Rigidity results

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i) If $\mu \leq \mu_{4}$, then $\Omega$ is a disk or an annulus.

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Assume that (Q) admits a solution. Then:
i) If $\mu \leq \mu_{4}$, then $\Omega$ is a disk or an annulus.
ii) If $\mu \leq \mu_{5}$ and $\partial \Omega$ has exactly two connected components, then $\Omega$ is an annulus.

The proof uses some of the ideas of Avilés and Deng, together with a comparison result between Dirichlet and Neumann eigenvalues (Friendlander '95, Filonov '05).
In case ii) we also need to make use of the result of Reichel for overdetermined problems in annuli.

## Main result

The answer to the question is NO !

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The answer to the question is NO!
Theorem
There exist nonradial domains $\Omega=\Omega_{0} \backslash \underline{\Omega_{1}}$, where $\Omega_{i} \subset \mathbb{R}^{2}$ are bounded smooth and simply connected, $\overline{\Omega_{1}} \subset \Omega_{0}$, such that the problem:

$$
\begin{cases}-\Delta u=\mu u & \text { in } \Omega \\ u=c_{i} & \text { on } \Gamma_{i} \\ \partial_{\nu} u=0 & \text { on } \partial \Omega\end{cases}
$$

admits a solution. Here $\mu>0$ and $\Gamma_{i}=\partial \Omega_{i}$.

## Connection to stationary Euler flows in 2D

The stationary Euler equations are:

$$
\begin{cases}v \cdot \nabla v+\nabla p=0 & \text { in } \mathbb{R}^{2}  \tag{E}\\ \operatorname{div} v=0 & \text { in } \mathbb{R}^{2}\end{cases}
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Here $v: \mathbb{R}^{2} \rightarrow \mathbb{R}^{2}$ is the velocity vector field and $p: \mathbb{R}^{2} \rightarrow \mathbb{R}$ is the pressure.

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Since $\operatorname{div} v=0$, we can define a stream function $\phi$ with $\nabla \phi=v^{\perp}$.
Any radial function $\phi$ gives rise to a solution of the Euler equations. In such case the trajectories (streamlines) are circles.
In particular, one can take compactly supported radial functions.
It is not known if there are compactly supported $C^{1}$ solutions to
(E) with noncircular streamlines.

## Connection to stationary Euler flows in 2D

Take $u$ as given by our main theorem. Then, define
$v=\left\{\begin{array}{ll}(\nabla u)^{\perp} & \text { in } \Omega, \\ 0 & \text { in } \mathbb{R}^{2} \backslash \Omega,\end{array} \quad p= \begin{cases}\mu\left(\frac{c_{0}^{2}}{2}-\frac{c_{1}^{2}}{2}\right) & \text { in } \overline{\Omega_{1}}, \\ -\frac{1}{2}\left(|\nabla u|^{2}+\mu\left(u^{2}-c_{0}^{2}\right)\right) & \text { in } \Omega, \\ 0 & \text { in } \mathbb{R}^{2} \backslash \Omega_{0} .\end{cases}\right.$
Then $(v, p)$ are continuous weak solution to (E) with non-circular streamlines and compact support.

## Connection to stationary Euler flows in 2D

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Then $(v, p)$ are continuous weak solution to (E) with non-circular streamlines and compact support.

Another compactly supported continuous weak solution of (E) has been given by Gómez-Serrano, Park and Shi. Their example is of vortex-patch type and the streamfunction does not solve any elliptic PDE.

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## Sketch of the proof

We consider a 1-parametric family of annuli $A(0 ; a, 1)$, where $a \in(0,1)$.
We denote $\mu_{0, k}(a)$ the $k$-th radial eigenvalue, with $\mu_{0,0}(a)=0$. We take $\psi_{a}$ the radial eigenfunction related to $\mu_{0,2}(a)$.

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Our aim is to show that, for some $a^{*} \in(0,1)$, there exists a branch of nonradial solutions of $(\mathrm{Q})$ bifurcating from $A\left(0 ; a^{*}, 1\right)$ and $\psi_{a^{*}}$.

We cannot perform this argument from $\mu_{0,1}(a)$ : in this case, the eigenfunction attains its minimum and maximum on the boundary. Hence also a bifurcating solution would do so. But in this case it is known that $\Omega$ has to be an annulus (Reichel '95).

## Sketch of the proof



The second radial Neumann eigenfunction in the annulus $A(0 ; a, 1)$

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## A direct approach

In polar coordinates $(r, \theta) \in \mathbb{R}^{+} \times \mathbb{S}^{1}$, we define:

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\Omega_{a}^{b, B}=\{(r, \theta): a+b(\theta)<r<1+B(\theta)\} .
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For small $(b, B) \in C^{k, \alpha}\left(\mathbb{S}^{1}\right)^{2}$, we can solve:

$$
\begin{cases}-\Delta u=\mu u & \text { in } \Omega_{a}^{b, B}, \\ \partial_{\nu} u=0 & \text { on } \partial \Omega_{a}^{b, B} .\end{cases}
$$

Then we can define $F(b, B)$ and $\tilde{F}(b, B) \in C^{k, \alpha}\left(\mathbb{S}^{1}\right)^{2}$,

$$
\begin{gathered}
F(b, B)(\theta)=(u(a+b(\theta), \theta)), u(1+B(\theta), \theta)), \\
\tilde{F}(b, B)(\theta)=\widetilde{F(b, B)}(\theta), \text { where } \widetilde{f}=f-\frac{1}{2 \pi} \int_{\mathbb{S}^{1}} f
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The zeroes of $\tilde{F}$ are the solutions to our problem. Clearly, $\tilde{F}(0,0)=0$ for any $a \in(0,1)$.

## A direct approach

It turns out that:

$$
D \tilde{F}_{(0,0)}(\hat{b}, \hat{B})(\theta)=\left(c_{1}(a) \Psi(a, \theta), c_{2}(a) \Psi(1, \theta)\right)
$$

where $c_{i}(a) \neq 0$ are constants and $\Psi=\Psi_{\hat{b}, \hat{B}}$ solves:

$$
\begin{cases}-\Delta \Psi=\mu \Psi & a<r<1 \\ \partial_{\nu} u=\hat{b}(\theta) & r=a \\ \partial_{\nu} u=\hat{B}(\theta) & r=1\end{cases}
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which becomes degenerate if $\mu_{0,2}(a)$ is a Dirichlet eigenvalue! However, it turns out that:

$$
D \tilde{F}_{(0,0)}: C^{k, \alpha}\left(\mathbb{S}^{1}\right)^{2} \rightarrow C^{k+1, \alpha}\left(\mathbb{S}^{1}\right)^{2}
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## Loss of derivatives

Then the cokernel of $D \tilde{F}_{(0,0)}$ is infinite dimensional and classical bifurcation theorems cannot be applied. This is a typical case of loss of derivatives.

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However, a new approach that avoids this problem has been recently found by Fall, Minlend and Weth.

## The functional framework

Denote by $\Omega_{1 / 2}=A(0 ; 1 / 2,1)$, and $\Phi_{a}^{b, B}$ a diffeomorphism:

$$
\Phi_{a}^{b, B}: \Omega_{1 / 2} \rightarrow \Omega_{a}^{b, B}, L_{a}^{b, B}=\left(\Phi_{a}^{b, B}\right)^{*} \Delta+\mu_{0,2}(a) I d
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$$

For $b=B=0$ we have the 2nd Neumann eigenfunction $\psi_{a}$, and $\bar{\psi}_{a}=\psi_{a} \circ \Phi_{a}^{0,0}$.
We also define the function spaces:

$$
\mathcal{X}^{k, \alpha}=\left\{u \in C^{k, \alpha}\left(\Omega_{1 / 2}\right): \partial_{r} u \in C^{k, \alpha}\left(\Omega_{1 / 2}\right)\right\}
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\mathcal{X}_{D}^{k, \alpha}=\left\{u \in \mathcal{X}^{k, \alpha}: u=0 \text { on } \partial \Omega_{1 / 2}\right\} \\
\mathcal{X}_{D N}^{k, \alpha}=\left\{u \in \mathcal{X}_{D}^{k, \alpha}: \partial_{r} u=0 \text { on } \partial \Omega_{1 / 2}\right\}
\end{gathered}
$$

## The nonlinear operator

We define the operator:

$$
G_{a}(v)=L_{a}^{b_{v}, B_{v}}\left[\bar{\psi}_{a}+w_{v}\right]
$$

where $v$ is in a neighborhood of 0 in $\mathcal{X}_{D}^{2, \alpha}$,

$$
\begin{aligned}
b_{v}(\theta) & =c_{1}(a) \partial_{r} v\left(\frac{1}{2}, \theta\right) \\
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and $w_{v} \in \mathcal{X}_{D N}^{2, \alpha}$ is defined as:

$$
w_{v}(r, \theta)=v(r, \theta)+\frac{\bar{\psi}_{a}^{\prime}(r)}{2(1-a)}\left[2(1-r) b_{v}(\theta)+(2 r-1) B_{v}(\theta)\right] .
$$

## The linearization

## Proposition

The map $G$ is defined in a neighborhood of 0 in $\mathcal{X}_{D}^{2, \alpha}$ and takes values in $\mathcal{Y}$, defined as:

$$
\mathcal{Y}=C^{1, \alpha}\left(\bar{\Omega}_{\frac{1}{2}}\right)+\mathcal{X}_{D}^{0, \alpha}
$$

endowed with the norm
$\|u\|_{\mathcal{Y}}=\inf \left\{\left\|u_{1}\right\|_{C^{1, \alpha}}+\left\|u_{2}\right\|_{\mathcal{X}^{0, \alpha}}: u_{1} \in C^{1, \alpha}, u_{2} \in \mathcal{X}_{D}^{0, \alpha}, u=u_{1}+u_{2}\right\}$.

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Moreover, $D G_{a}(0)(v)=L_{a}^{0,0}(v)$ is a Fredholm operator of index 0 from $\mathcal{X}_{D}^{2, \alpha}$ to $\mathcal{Y}$.
This operator becomes degenerate if $\mu$ is a Dirichlet eigenvalue.

## Eigenvalues' crossing

Recall that $\mu_{0,2}(a)$ is the second Neumann eigenvalue which is radially symmetric in $A(0 ; a, 1)$.
We now denote $\lambda_{l, 0}(a)$ the first Dirichlet eigenvalue for functions of mode $I$, that is, $\psi(r, \theta)=\phi(r) \cos (I \theta)$.

## Proposition

The following asymptotics hold:

1. If $a$ is close to 1 , then $\mu_{0,2}(a)>\lambda_{l, 0}(a)$.
2. If $a$ is close to 0 and $I \geq 4$, then $\mu_{0,2}(a)<\lambda_{I, 0}(a)$.

Then, there exists some $a^{*} \in(0,1)$ for which $\mu_{0,2}\left(a^{*}\right)=\lambda_{l, 0}\left(a^{*}\right)$.

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Then, there exists some $a^{*} \in(0,1)$ for which $\mu_{0,2}\left(a^{*}\right)=\lambda_{l, 0}\left(a^{*}\right)$.
To exclude resonances with lower modes we restrict ourselves to $l$-symmetric functions, $I \geq 4$. Moreover this allows us to rule out the invariance by translations.

## Bifurcation

The degeneracy of $D G_{a}(0)$ is necessary to obtain bifurcation.


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By the classical Crandall-Rabinowitz theorem, bifurcation occurs if the transversality condition holds, that here reduces to:

$$
\mu_{0,2}^{\prime}\left(a^{*}\right)>\lambda_{1,0}^{\prime}\left(a^{*}\right)
$$

We have been able to prove this only for large $I$.

## Theorem

Let $I \geq I_{0}$ sufficiently large. There exist some $\varepsilon>0$ and a continuously differentiable curve

$$
\begin{gathered}
(-\varepsilon, \varepsilon) \ni s \mapsto\left\{\left(a(s), b_{s}, B_{s}\right) \subset(0,1) \times C^{2, \alpha}\left(\mathbb{S}^{1}\right) \times C^{2, \alpha}\left(\mathbb{S}^{1}\right),\right. \\
\left.\left(a(0), b_{0}, B_{0}\right)=\left(a^{*}, 0,0\right)\right\}, \text { and } \\
b_{s}(\theta)=s c \cos (I \theta)+o(s), B_{s}(\theta)=s C \cos (\mid \theta)+o(s), C>0>c,
\end{gathered}
$$

such that the following problem admits a solution:

$$
\begin{cases}\Delta u_{s}+\mu_{0,2}(a(s)) u_{s}=0 & \text { in } \Omega_{a(s)}^{b_{s}, B_{s}}, \\ \nabla u_{s}=0 & \text { on } \partial \Omega_{a(s)}^{b_{s}, B_{s}} .\end{cases}
$$

A posteriori, the domain and the solution are analytic.

## Numerical approximation

For specific values of $I$ one can give numerical approximations by using Mathematica. For instance, if $I=4$,

$$
\begin{gathered}
a^{*}=0,140989 \ldots, \quad \mu_{0,2}\left(a^{*}\right)=\lambda_{4,0}\left(a^{*}\right)=57,5851 \cdots=\mu_{18}, \\
\mu_{0,2}^{\prime}\left(a^{*}\right)=105,971 \ldots, \lambda_{4,0}^{\prime}\left(a^{*}\right)=0,12067 \ldots
\end{gathered}
$$



## Thank you for your attention!

